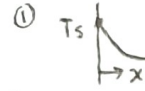
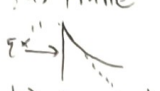


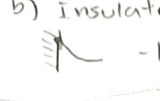
ME 109 midterm 1

conduction: $q'' [W/m^2] = -k \frac{dT}{dx}$ $k [W/m \cdot K]$
 convection: $q'' [W/m^2] = h(T_s - T_\infty)$ $h [W/m^2 \cdot K]$
 radiation: $q'' = \epsilon \sigma (T_s^4 - T_{sur}^4)$ $\sigma = 5.67 \times 10^{-8} [W/m^2 \cdot K^4]$
 OR $q [W] = hrA(T_s - T_{sur})$ $\epsilon = (0, 1)$ "emissivity"
 where $hr = \epsilon \sigma (T_s + T_{sur})(T_s^2 + T_{sur}^2)$ perfect reflector perfect absorber

Fourier's Law: $q_x = -KA \frac{dT}{dx}$ (1-D) [W]
 $\vec{q}'' = -k \nabla T = -k \left(\frac{\partial T}{\partial x} \hat{i} + \frac{\partial T}{\partial y} \hat{j} + \frac{\partial T}{\partial z} \hat{k} \right)$ [W/m^2]
 (3-D version, not area normalized)
 $\frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left(k \frac{\partial T}{\partial z} \right) + \dot{q} = \rho c_p \frac{\partial T}{\partial t}$
 $q''_{in} - q''_{out}$ Egen Est

Types of Boundary conditions

- ① const. surf. Temp. $T(0, t) = T_s \forall t$

- ② constant heat fluxes.
 a) finite $-k \frac{\partial T}{\partial x} |_{x=0} = q''_0$

- ③ Convection. $-k \frac{\partial T}{\partial x} |_{x=0} = h(T_\infty - T(0, t))$

- b) insulated (adiabatic) $-k \frac{\partial T}{\partial x} |_{x=0} = 0$


Variable Thermal Conductivity ($K = K_0(1 + BT)$)

$-K \frac{\partial^2 T}{\partial x^2} = q''$ (const for GEOMETRICALLY symmetric & no external heat gen $\dot{q}(x)$)
 $\rightarrow B=0 \quad q'' = -\frac{K_0}{L} [(T_2 - T_1) + \frac{B}{2} (T_2^2 - T_1^2)]$
 // solved via sep of variables

Finns 
 $q_x - q_{x+dx} = dq_{conv} = h(P dx)(T(x) - T_\infty)$
 $q_x = KA(x) \frac{dT}{dx}$
 Gov Eq: $\frac{d}{dx} \left(-KA(x) \frac{dT}{dx} \right) - hP(T(x) - T_\infty) = 0$
 $\Rightarrow \frac{d^2 T}{dx^2} - \frac{hP}{KAC} [T(x) - T_\infty] = 0 \Rightarrow \frac{d^2 \theta}{dx^2} - \frac{hP}{KAC} \theta = 0$

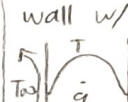
Aside: $\frac{\partial^2 \theta}{\partial x^2} + m^2 \theta = 0 \Rightarrow C_1 e^{mx} + C_2 e^{-mx}$ (homogeneous 2nd order DE!)
 Key substitutions: $m^2 = hP/KAC$ $\theta = T(x) - T_\infty$

Fin effectiveness & efficiency

$\epsilon_s = \frac{q_s}{hA\theta_b} \Rightarrow \frac{\text{fin dissipation}}{\text{bare wall diss.}}$ $\frac{10t}{t} \Rightarrow \frac{10t}{t} = \frac{20}{1}$
 For infinite fin, $\epsilon_s = \sqrt{\frac{KP}{hAC}}$ Note: $\epsilon_s \uparrow \leftrightarrow \frac{P}{AC} \uparrow$
 $\eta_f = q_t/q_{max} = q_t/hA_f \theta_b$ A_f : total Area!

Spheres & cylinders


 $q_r = T_1 - T_2 / \left(\frac{\ln(r_o/r_i)}{2\pi KL} \right)$ [W]
 $T(r) = T_1 - \frac{T_1 - T_2}{\ln(r_o/r_i)} \ln(r/r_i)$
 Assume $L \gg r_i, r_o$
 $q = T_1 - T_2 / R_{sph}$ where $R_{sph} = \frac{r_o - r_i}{4r_o r_i \pi k}$
 For solid sphere w/ heat gen. $T(r) = -\frac{\dot{q} r^2}{6k} + \frac{\dot{q} r_o}{3h} + \frac{\dot{q} r_o^2}{6k} + T_\infty$

wall w/ heat gen 
 Gov Eq: $\frac{\partial^2 T}{\partial x^2} = -\frac{\dot{q}}{k}$
 quick math $\Rightarrow T(x) = -\frac{\dot{q}}{2k} x^2 + C_1 x + C_2$

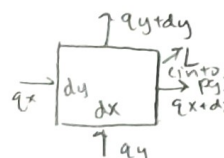
Infinite fin example.

lim $T(x) = C_1 e^{mx} + C_2 e^{-mx} \rightarrow 0 = C_1 e^{mx}$
 $\theta(x) = \theta_b e^{-mx}$
 $q_s = hPAC \theta_{bound} = -KAC \frac{d\theta}{dx} |_{x=0}$

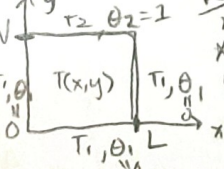
Systems of Fins


 $A_{total} = A_{bare wall} + N \cdot A_{fin}$
 $q_{tot} = N \eta_f A_f \theta_b h + h \theta_b A_b$
 $q_{tot} = h A_{tot} \left[1 - \frac{NA_f}{A_{tot}} (1 - \eta_f) \right] \theta_b$
 $\eta_o = q_{tot} / h A_{tot} \theta_b = \left[1 - \frac{NA_f}{A_{tot}} (1 - \eta_f) \right]$

2-D Steady Heat Conduction


 $q_x + q_y = q_{x+dx} + q_{y+dy}$ // energy conservation.
 $\Rightarrow$ Taylor expand to first term.
 $q_x = -K(dyL) \frac{\partial T}{\partial x}$
 $q_y = -K(dxL) \frac{\partial T}{\partial y}$
 Gov Eq: $\frac{\partial}{\partial x} \left(K \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(K \frac{\partial T}{\partial y} \right) = 0$
 // assuming K is const
 $\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0 \Rightarrow \nabla^2 T = 0$

standard basis for Superposition

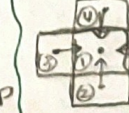

 Sub: $\theta = \frac{T - T_1}{T_2 - T_1}$
 $x=0, T=T_1, \theta=0$
 $x=L, T=T_1, \theta=0$
 $y=0, T=T_1, \theta=0$
 $y=W, T=T_2, \theta=1$
 Sol: $\theta(x, y) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n} \sin \frac{n\pi x}{L} \frac{\sinh(n\pi y/L)}{\sinh(n\pi W/L)}$

Shape Factors

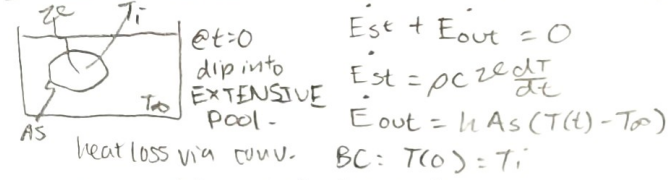
- ① isothermal sphere buried in semi-inf med.
- ② horz cylinder L buried in semi-inf med.
- ③ vert cylinder in semi-inf med.
- ④ Between 2 cylinders in infinite medium.

Restriction $z \gg D/2$ $S = \frac{2\pi D}{1 - D/4z}$ $T_1 - T_2$
 $L \gg D$ $S = \frac{2\pi L}{\ln(4z/D)}$ T_2
 $z \gg 3D/2$ $S = \frac{2\pi L}{\ln(4L/D)}$ no z!
 $L \gg D_1, D_2$ S' = some shit.
 $L \gg W$
 $q = S k \Delta T$ $R_{th} = \frac{1}{S k}$ $S = [L]$

NODAL analysis


 $\dot{q}_0 + \sum_k \dot{q}_k = 0$ // energy balance.
 shortcut: when $\Delta x = \Delta y$ & No heat gen. $T_{m,n} = \frac{T_2 + T_3 + T_4 + T_5}{4}$
 Nodal with conv & with composite materials. ---

Transient Heat conduction (No Spatial Temp gradient)



Gov. Eq: $\frac{dT}{dt} = -\frac{hA_s(T-T_\infty)}{\rho C V}$ $T(0) = T_i$
 let $\theta = T - T_\infty$ solve.
 $T(t) = (T_i - T_\infty) e^{-\frac{hA_s}{\rho C V} t} = (T_i - T_\infty) e^{-t/\tau}$
 where $\tau = \left(\frac{\rho C V}{hA_s}\right)^{-1}$

Lumped Capacitance model applies when $Bi < 0.1$

Nonhomogenous Lumped Capacitance

Plane wall w/ convection. Gov Eq: $\frac{dT}{dt} + \frac{h}{\rho C \delta} (T - T_\infty) = \frac{q''}{\rho C \delta}$
 solve $\theta(t) = \left(-\frac{q''}{h}\right) e^{-\frac{h}{\rho C \delta} t} + \frac{q''}{h}$
 Math aside: $\frac{d\theta}{dt} + m\theta = C$
 $\Rightarrow \theta(t) = C_1 e^{-mt} + \frac{C}{m}$ C_1 determined via BC.

Cylinder w/ convection. $\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \frac{\partial T}{\partial r} = \frac{1}{\alpha} \frac{\partial T}{\partial t}$
 $\theta^* = \sum_{n=1}^{\infty} C_n \exp(-\zeta_n^2 Fo) J_0(\zeta_n r^*)$
 Now, $Fo = \alpha t / r_o^2$ $Bi = \frac{h r_o}{k}$
 $\frac{Q}{Q_0} = 1 - \frac{2\theta_0^*}{\zeta_1} J_1(\zeta_1)$

Sphere w/ convection. $\frac{\partial^2 T}{\partial r^2} + \frac{2}{r} \frac{\partial T}{\partial r} = \frac{1}{\alpha} \frac{\partial T}{\partial t}$
 $\theta^* = \sum_{n=1}^{\infty} C_n \exp(-\zeta_n^2 Fo) \frac{1}{\zeta_n r^*} \sin(\zeta_n r^*)$
 $\frac{Q}{Q_0} = 1 - \frac{3\theta_0^*}{\zeta_1} [\sin(\zeta_1) - \zeta_1 \cos(\zeta_1)]$

Discrete Transient Numerical Analysis

Gov Eq: $K A \frac{T_{i-1}^p - T_i^p}{\Delta x} - K A \frac{T_i^p - T_{i+1}^p}{\Delta x} = \rho C A \Delta x \frac{T_i^{p+1} - T_i^p}{\Delta t}$
 $T_i^{p+1} = Fod [T_{i-1}^p + T_{i+1}^p] + (1 - 2Fod) T_i^p$
 $Fod = \frac{\Delta t}{\alpha \Delta x^2}$
 $Bi_d = \frac{h \Delta x}{k}$ discrete versions coefficient > 0 for STABILITY

Nonhomogenous 1-D conduction problem

Energy balance setup: $dq_{conv} + dq_{gen} = dq_{cond}$
 $K A_c \frac{d^2 T}{dx^2} dx + \dot{q} A_c dx = h P dx (T - T_\infty)$
 $\Rightarrow \frac{d^2 T}{dx^2} - \frac{hP}{K A_c} (T - T_\infty) = -\dot{q}/K$
 Sub $\theta = T - T_\infty$
 $\theta_c = C_1 e^{mx} + C_2 e^{-mx}$
 $\theta_p = \text{some const}$
 critical def $\frac{d^2 \theta_p}{dx^2} = 0$ $-m\theta_p = -\dot{q}/K \Rightarrow \theta_p = \dot{q}/mK$
 $\theta(x) = C_1 e^{mx} + C_2 e^{-mx} + \dot{q}/mK$

Critical Dimensional Quantities

$Bi = \frac{h L_c}{k}$ or $\frac{h r_o}{k}$ for conv/rad resp.
 // ratio of conv to conduction capability
 $Fo = \alpha t / L_c^2$ // dimensionless time. $\alpha = \frac{k}{\rho C}$ (thermal diffusivity)
 $\frac{L_c}{As} = L_c$ // characteristic length (relevant in Bi, Fo)

Spatial effects when $Bi > 0.1$.

Dimensionless Substitutions: Gov Eq: $\theta^* = \frac{\theta}{\theta_i} = \frac{T - T_\infty}{T_i - T_\infty} = f(x^*, Fo, Bi)$
 $x^* = x/L$ $t^* = \alpha t / L^2 = Fo$
 $\frac{\partial^2 \theta^*}{\partial x^{*2}} = \frac{\partial \theta^*}{\partial t^*}$

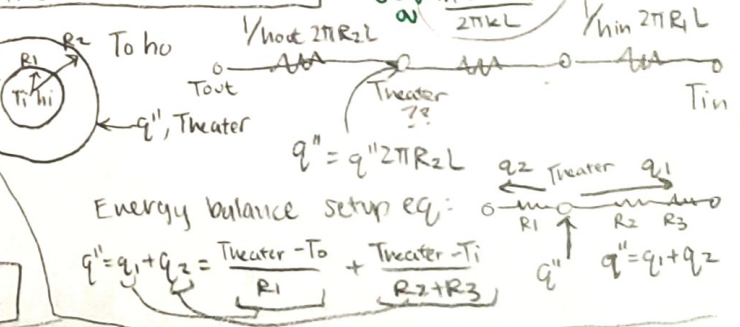
Plane wall w/ convection $\frac{\partial^2 T}{\partial x^2} = \frac{1}{\alpha} \frac{\partial T}{\partial t}$
 $\theta^* = \sum_{n=1}^{\infty} C_n \exp(-\zeta_n^2 Fo) \cos(\zeta_n x^*)$ where C_n, ζ_n can be found on table.
 // When $Fo > 0.2$, can use a first-term approximation.

$\frac{Q}{Q_0} = 1 - \frac{\sin \zeta_1}{\zeta_1} \theta_0^*$ where $\theta_0^* = \frac{T(x=0,t) - T_\infty}{T_i - T_\infty}$ // temperature @ centerline
 ratio of total e transferred / max e transferrable

Semi-Infinite Solids

Case 1: Const. Surface BC Temp $T(0,t) = T_s$
 $\frac{T(x,t) - T_s}{T_i - T_s} = \text{erf}\left(\frac{x}{2\sqrt{\alpha t}}\right)$
 Case 2: Const. Surface heat flux: $q_s'' = q_0''$
 $T(x,t) - T_i = \frac{2q_0''(\alpha t / \pi)^{1/2}}{k} \exp\left(-\frac{x^2}{4\alpha t}\right) - \frac{q_0'' x}{k} \text{erfc}\left(\frac{x}{2\sqrt{\alpha t}}\right)$
 Case 3: Surface convection $-k \frac{\partial T}{\partial x}|_{x=0} = h [T_\infty - T(0,t)]$
 $\frac{T(x,t) - T_i}{T_\infty - T_i} = \text{erfc}\left(\frac{x}{2\sqrt{\alpha t}}\right) - \left[\exp\left(\frac{hx}{k} + \frac{h^2 \alpha t}{k^2}\right)\right] \left[\text{erfc}\left(\frac{x}{2\sqrt{\alpha t}} + \frac{h\sqrt{\alpha t}}{k}\right)\right]$

Practice Midterm P#1



ME109 Midterm 2. Cheatsheet.

Boundary Layers δ = velocity BL δ_T = Thermal BL.

Local Surface heat flux: $q_s'' = -k_f \frac{\partial T}{\partial y} |_{y=0}$

Newton's Law of Cooling: $q_s'' = h(T_s - T_\infty)$

Local heat trans. coeff

Fourier's Law?

$$h = \frac{-k_f \frac{\partial T}{\partial y} |_{y=0}}{(T_s - T_\infty)}$$

$\bar{h} \triangleq \frac{1}{L} \int_0^L h_x dx \rightarrow \bar{q} = \bar{h} A_s (T_s - T_\infty)$ // general heat transfer of ENTIRE plate.

Velocity BL

$$C_f = \frac{\tau_s}{\rho u_\infty^2 / 2}$$

$$\tau_s = \mu \frac{\partial u}{\partial y} |_{y=0}$$

Thermal BL

$$\delta_T \triangleq \frac{T_s - T(\delta_T)}{T_s - T_\infty} = 0.99$$

$$\delta \triangleq \frac{u}{u_\infty} = 0.99$$

1. IMPORTANT DIM-LESS Quantities

& other relations.

$$Re_x = \frac{u_\infty x}{\nu}$$

$$\nu = \mu / \rho$$

$$\alpha = \frac{k}{\rho c_p}$$

→ internal flow $Re_{cr} = 2300$

→ external flow $Re_{cr} = 5 \times 10^5$

$$Pr = \frac{c_p \mu}{k_f} = \frac{\nu}{\alpha} \approx \frac{\delta}{\delta_T}$$

Relative BL thickness $\sim Pr$!

$$Nu_x = \frac{h_x L}{k_f}$$

$$\bar{Nu}_L = \frac{\bar{h} L}{k_f}$$

ALWAYS TRUE REGARDLESS OF TURBULENCE.

$$Ra_x = Gr_x Pr = \frac{g \beta (T_s - T_\infty) x^3}{\nu \alpha}$$

$$Gr_x = \frac{g \beta (T_s - T_\infty) x^3}{\nu^2}$$
 Grashof #

Sphere:

$$\bar{Nu}_D = 2 + (0.4 Re_D^{1/2} + 0.06 Re_D^{2/3}) Pr^{0.4} \left(\frac{\mu}{\mu_s} \right)^{1/4}$$

$$\bar{Nu}_D = C Re_D^m Pr^{1/3}$$

→ C & m from T7.2

// most common conditions $0.4 \leq Re_D \leq 4 \times 10^5$
 $Pr \geq 0.7$

Fluid Mechanics Gov Eq Review

Mass Cons.

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

Momentum

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2}$$

Energy: $u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha \frac{\partial^2 T}{\partial y^2}$

External Flows

Laminar Flat Plate + exact integral method.

Turbulent Flat Plate * discussion printout.

Mixed flat plate

Cylinders

$$Re_D = \frac{u_\infty D}{\nu}$$

$$\bar{Nu}_D = \frac{\bar{h}_D D}{k_f}$$

Internal Flows

$$Deff = \frac{4 A_c}{Perim.}$$

// noncircular cylinder

circular Laminar

$$Nu_D = 4.36$$

unif. q_s'' } fully developed

$$Nu_D = 3.66$$

unif T_s

$$\bar{Nu}_D = 3.66 + \frac{0.0668 (D/L) Re_D Pr}{1 + 0.04 [(D/L) Re_D Pr]^{2/3}}$$

// thermal entry, $Pr \geq 5$

Turbulent.

$$Nu_D = 0.023 Re_D^{4/5} Pr^{1/3}$$

// most common

$Re \geq 10,000$

$0.6 \leq Pr \leq 160$

$L/D \geq 10$

$$n=0.4$$

OR

$$n=0.3$$

Thermal vs Hydro entry lengths

$$L_{sd,t} = 0.05 Re_D Pr$$
 thermal

hydro

$$L_{sd,h} = 0.05 Re_D$$

Internal Flows Temp. Relations

const q_s''

$$\frac{dT_m}{dx} = \frac{q_s'' P}{\dot{m} c_p} \neq f(x) \quad !$$

$$T_m(x) = T_{mi} + \frac{q_s'' P}{\dot{m} c_p} x \quad \star$$

$q_s'' = \text{const}$

const T_s

$$\frac{dT_m}{dx} = \frac{P}{\dot{m} c_p} h \Delta T$$

$$\frac{T_s - T_m(x)}{T_s - T_{mi}} = \exp\left(-\frac{Px}{\dot{m} c_p h}\right)$$

$$\frac{T_s - T_{mo}}{T_s - T_{mi}} = \frac{\Delta T_o}{\Delta T_i}$$

Additional

$$T_m = \frac{\int_{Ac} \rho u(r) c_p T(r) dAc}{\dot{m} c_p}$$

// Bulk mean temp.

$$\Delta T_m = (\Delta T_i + \Delta T_o)/2 \quad // \text{mean } \Delta T$$

$$\Delta T_{lm} = \frac{\Delta T_o - \Delta T_i}{\ln(\Delta T_o / \Delta T_i)} \quad // \text{log-mean } \Delta T$$

Natural Convection



vertical plates

$$Nu_L = C Ra_L^n$$

Lam: $C = 0.59 \quad n = \frac{1}{4}$

Turb: $C = 0.10 \quad n = \frac{1}{3}$

For Laminar: $Ra \leq 10^9$

Horz Plates



$$Nu_L = 0.54 Ra_L^{1/4}$$

$$Ra_L \leq 10^7$$

$$Nu_L = 0.15 Ra_L^{1/3}$$

$$Ra_L \geq 10^7$$



$$Nu_L = 0.27 Ra_L^{1/4}$$



$$Nu_D = C Ra_D^n$$

Table 9.1 for C, n

Practice Final

$$\frac{T_{\infty} - T_{mo}}{T_{\infty} - T_{mi}} = \exp\left(-\frac{UA}{\dot{m} c_p}\right)$$

Temperature difference in convective tube

